

Research Article

MHD Jeffrey Hybrid NANOFLUID Flow with Vanadium Pentoxide (V_2O_5) and Thermal Radiation Squeezed Between Two Parallel Plates

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Received: 25/Feb/2024; Accepted: 27/Mar/2024; Published: 30/Apr/2024

Abstract— this work exposes the properties of the transmission of heat and mass in Jefferies nanofluids that are hybrid over the squeezing plates that travel across a porous material, as well as the nature of magnetohydrodynamics (MHD) fluid flow. The influence of thermal radiation parameters together with other pertinent parameters are studied. Suitable similarity transformation is applied to the system of dimensionless equations to discretize them. Vanadium Pentoxide (V_2O_5) dispersions are deliberated in the base fluid. Validation of this research is achieved by relating to published results. Graphs and tables discuss the momentum, and temperature, with Nusselt number and concentration profiles. The graphical results show that the temperature profile elevates with $S < 0$, $S > 0$, Ha , Ec and Rd , while it drops for De and λ . The Nusselt number profile had increased with an increase in Pr , while it decreased with an increase in the thermal radiation parameter Rd . While Ha and λ increase fluid flow concentration while lowering the mass exchange rate, they have a major negative effect on the thermal radiation parameter Rd .

Keywords— Jeffrey hybrid nanofluid, MHD, Squeeze flow, Thermal radiation, Chemical reaction

1. Introduction

Convective heat transfer in nanofluids has drawn a lot of attention in the last few years. Choi [1] first used the term "nanofluid" to refer to a liquid solution that contains tiny particles (diameter less than 50 nm). According to experimental investigations (e.g., Madaki et al., [2] and [3]), the base liquid's thermal conductivity can be improved by 10-50%, with a notable improvement in the convective heat transfer coefficient, even with a modest volumetric proportion of nanoparticles (typically less than 5%). The volatile flow in two dimensions of a viscous magnetohydrodynamic MHD fluid across two parallel inconceivable plates was investigated and analysed by Akbar et al. [4]. A review of the literature reveals that the squeeze Jeffrey flow across two parallel plates is the subject of most studies. Squeezing flow is the term used to describe the constriction of fluid across two separate plates that results in the fluid beginning to flow. Because it is used in the geometric model of lubricant flow in pneumatic elevators, moulding with injections, and bearings as well, it has been researched by several scientists, such as Madaki et al. [5], [6], [7], Hussaini et al. [8], [9]. Wang [10] discovered an entirely novel similarities transformation for Navier-Stokes into ODEs viscous fluid of squeezing flow. Bujurke et al. [11] further examined a logical clarification. Rashidi et al. [12], along with Khan et al. [13]. Scientists are quite interested in Jeffrey fluid because it exhibits both relaxation and retardation characteristics. How the fluid

behaves in response to applied shear stress. Because it can change from a shear thinning flow to a Newtonian fluid when a high force is applied to the fluid, it is categorized as such. Noor et al. [14]. Matinfar et al. [15] went into great detail about the impact of thermal radiation on squeezing, Jeffery fluid fluxes. The impact of thermal radiation and heat generation/absorption on the squeezing of unstable cu as well as TiO_2 -nanofluid was investigated quantitatively as well as computationally by Madaki et al. [16]. Mamatha et al. [17] investigated the effect of electromagnetic radiation on non-steady free convection magnetohydrodynamic flows of Brinkman-type systems over a porous media that is heated by Newtonian radiation. Thermal radiation's influence on squeezing flow fluid from Casson between parallel plates was reported by Khan et al. [18]. Makinde et al. [19] examined the effects of Brownian motion and thermophoresis on MHD bioconvection of nanofluid through a quartic chemical process along with the nonlinear thermal radiation beyond an upper horizontal surface of a paraboloid of rotation. The influence of an integrated electromagnetic field on radiant bioconvection flow via a plate that is vertical having thermophoresis along with Brownian motion was studied by Avinash et al. [20]. The impacts of suspended nanoparticles and non-linear thermal radiation on the convection and heat transfer boundary layer flow of nanofluids on a heated vertical sheet were examined by Mahanthesh et al. [21]. The purpose of heat generation/absorption is to decrease and increase a fluid's thermal conductivity, correspondingly. The

temperature of the high-conductivity fluid rises, while the low-conductivity fluid exhibits the reverse behaviour. Noor et al. [22]. Using the power series technique, Qasim [23] examined the effects of the heat sink/source around the temperature and mass movement on a vertical stretching plate.

Many studies have demonstrated the use of hybrid nanofluid in high-tech applications, including coolant in automobile engines, nuclear power plants, and electrical appliances (Ahmed et al., [24]). Hybrid nanofluid flow has been demonstrated in several physical models. Alghamdi et al. [25] investigated the MHD flow and heat transmission for the Casson hybrid nanofluid problem. After that, Abbas et al. [26] looked into how stagnation point flow at vertical stretched surfaces was affected by viscous dissipation, MHD, thermal slip, and nonlinear radiative heat transfer. Mahabaleshwar et al. [27] investigated the MHD flow and radiative heat transfer of a Cu–Al2O3/H2O hybrid nanofluid placed under velocity slip conditions over a porous stretch/shrink surface.

There has not yet been any research done on the mass along with heat transmission of the MHD composite nanofluid through the squeezing of a pair of parallel plates acting as thermal radiation sources, with vanadium pentoxide (V2O5) acting as the solid material (nanoparticle), to the best of our knowledge. Squeezing Jeffrey's hybrid nanofluid flow due to its complexity in terms of governing equations has generally received little attention. Therefore, this research is aimed at obtaining the mathematical solutions of the Mass and heat transfer in radiant-MHD vanadium pentoxide (V2O5)-based squeezing flow Jeffrey nanofluid hybrids are equipped with a source or heat generation/ absorption (i.e., filling the void left by Noor and Shafie [28]), and hence compare the present result with the existing ones.

2. Description of the Problem

This research will contain a numerical solution describing the behaviour of the hybrid nanofluid Jeffrey squeezed between a pair of parallel plates. However, our study would focus on the effects of radiant heating, magnetic field on the fluid flow and the heat source/sink on the general flow pattern. Thus, the fluid flow would consider to be squeezing, Newtonian and incompressible.

We study the unstable squeezed Jeffrey hybrid nanofluid flow across a pair of plates over a porous material with magnetohydrodynamic, heat generation/absorption as well as chemical reaction. Vanadium peroxide hybrid nanoparticles are taken into consideration. There is a detachment y between two sheets as shown on the figure1 below. Both the top and bottom plates are subject to the peripheral speed of v_2 also as shown on fig1 below. Up to $t=1/\alpha$, the two separate sheets shift away when $\alpha < 0$ but nearer when $\alpha > 0$, just as shown on fig. 1 below. The bottom surface is exerted with magnetic field, $B(t) = B_0(1 - \alpha t)^{-1/2}$ vertically Noor et al. [29].

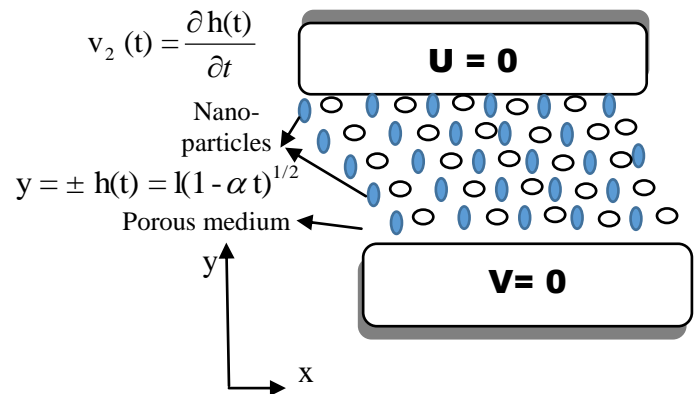


Figure 1. The fluid flow model.

Muhammad et al. ([30], [31]) articles contain a more general formulation of the Jeffrey hybrid nanofluid's momentum, temperature, nanoparticle concentration, and continuity equations.

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \tag{1}$$

$$\begin{aligned} \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} &= \frac{\mu_{hnf}}{\rho_{hnf}} \left(1 + \frac{1}{\lambda_1} \right) \frac{\partial^2 v}{\partial y^2} + \frac{\mu_{hnf}}{\rho_{hnf}} \frac{\lambda_2}{1 + \lambda_1} \\ &\left(\frac{\partial^3 v}{\partial t \partial y^2} + u \frac{\partial^3 u}{\partial x \partial y^2} + v \frac{\partial^3 u}{\partial y^3} - \frac{\partial u}{\partial x} \frac{\partial^2 u}{\partial y^2} + \frac{\partial u}{\partial y} \frac{\partial^2 u}{\partial x \partial y} \right) \\ - \frac{\sigma_{hnf} B(t)}{\rho_{hnf}} u - \frac{\mu_{hnf}}{\rho_{hnf}} \left(1 + \frac{1}{\lambda_1} \right) \frac{\phi}{k_1(t)} u &= 0, \end{aligned} \tag{2}$$

$$\begin{aligned} \frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} &= \frac{k_{hnf}}{(\rho C p)_{hnf}} \frac{\partial^2 T}{\partial y^2} + \frac{\mu_{hnf}}{(\rho C p)_{hnf}} \left(1 + \frac{1}{\lambda_1} \right) \left[4 \left(\frac{\partial u}{\partial x} \right)^2 + \left(\frac{\partial u}{\partial y} \right)^2 \right] \\ - \frac{1}{(\rho C p)_{hnf}} \frac{\partial q_r}{\partial y} + \frac{Q(T - T_\infty)}{\rho_{hnf}} &= 0, \end{aligned} \tag{3}$$

$$\frac{\partial C}{\partial t} + u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D_m \frac{\partial^2 C}{\partial y^2} - k_c (C - C_\infty) = 0, \tag{4}$$

Subject to the boundary conditions

$$u = 0, v = v_w = \frac{\partial h(t)}{\partial t}, T = T_w, C = C_\infty, \text{ at } y = h(t), \tag{5}$$

$$\frac{\partial u}{\partial y} = \frac{\partial^3 u}{\partial y^3} = \frac{\partial T}{\partial y} = \frac{\partial C}{\partial y} = v_0, \text{ at } y = 0, \tag{6}$$

Table 1.0 Thermophysical properties of the fluid and solid particles, Madaki et al. [32], [33]

Physical Properties	Fluid phase (water)	Vanadium Pentoxide (V ₂ O ₅)
C_p (J/Kg k)	4179	127.7
ρ (Kg/m ³)	997.1	3.35
K (W/mK)	0.613	4.22
σ ($\mu S/cm$)	0.05	13.0

According to the Stefan-Boltzmann rule, the amount of radiation absorbed q_r over period of time departing an object corresponds with a fourth of a degree of the temperature in absolute terms, or $q_r = \sigma T^4 A$. This represents the radiation heat fluctuation. Where T is the fundamental temperature, A is the emitting body's area, σ is the Stefan Boltzmann constant, and q is the heat transfer per unit time (ω). Nevertheless, Roseland [34] also clearly expressed the radiative heat flux factor that appears in equation (4) as

$$q_r = -\frac{4\sigma^* \partial T^4}{3k^* \partial y}$$

According to some studies by [4] and [35], it is thought that the term T^4 may be represented as a function of temperature and that the temperature differential inside the flow is significantly limited. As a result, the higher-order constants are ignored and T^4 is enlarged using the Taylor series expansion around T_∞ . As a result, we devise

$$T^4 = T_\infty^4 + 4T_\infty^3 T - 4T_\infty^2 T^2 + 4T_\infty T^3 - 3T_\infty^4$$

Substituting eqs. (7) and (8) into (3), gives

$$\frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \frac{k_{hnf}}{(\rho Cp)_{hnf}} \frac{\partial^2 T}{\partial y^2} + \frac{\mu_{hnf}}{(\rho Cp)_{hnf}} \left(1 + \frac{1}{\lambda_1}\right) \left[4 \left(\frac{\partial u}{\partial x}\right)^2 + \left(\frac{\partial u}{\partial y}\right)^2\right] + \frac{16\sigma^* T_\infty^3}{3k^*} \frac{\partial^2 T}{\partial y^2} + \frac{Q(T - T_\infty)}{\rho_{hnf}} = 0, \tag{9}$$

3. Method of solution with Runge-Kutta fourth order and shooting scheme combined

We used the same non-dimensional variables as in [28] to transform the mechanism of partial differential system (PDEs) into ordinary differential system (ODEs).

$$u = \frac{ax}{2(1-at)} f'(\eta), v = -\frac{al}{2\sqrt{(1-at)}} f(\eta), \tag{10}$$

$$\eta = \frac{y}{l\sqrt{(1-at)}} f'(\eta), \theta = \frac{T}{T_w}, \varphi = \frac{C}{C_w}$$

When the similarity variables (8) are substituted into equations (1)–(4), the resulting non-dimensional ODE is:

$$\frac{\mu_{hnf}}{\mu_f} \frac{\rho_f}{\rho_{hnf}} \left(1 + \frac{1}{\lambda_1}\right) f^{iv} - S(\eta f''' + 3f'' + ff'' - ff''') + \frac{\mu_{hnf}}{\mu_f} \frac{\rho_f}{\rho_{hnf}} \left(1 + \frac{1}{\lambda_1}\right) \frac{De}{2} (\eta f^{fv} + 5f^{iv} + 2ff'' - ff^{iv} - ff^{fv}) - \frac{\sigma_{hnf}}{\sigma_f} \frac{\rho_f}{\rho_{hnf}} Ha^2 f'' - \frac{\mu_{hnf}}{\mu_f} \frac{\rho_f}{\rho_{hnf}} \left(1 + \frac{1}{\lambda_1}\right) \frac{1}{Da} f'' = 0, \tag{11}$$

$$\frac{(\rho Cp)_f}{(\rho Cp)_{hnf}} \frac{k_{hnf}}{k_f} \frac{1}{Pr} \left(1 + \frac{4}{3} Rd\right) \theta'' + S(f\theta' - \eta\theta') + \frac{(\rho Cp)_f}{(\rho Cp)_{hnf}} \gamma \theta + \frac{\mu_{hnf}}{\mu_f} \frac{(\rho Cp)_f}{(\rho Cp)_{hnf}} Ec \left[\left(1 + \frac{1}{\lambda_1}\right) [(f'')^2 + 4\delta^2 (f')^2]\right] = 0, \tag{12}$$

$$\frac{1}{Sc} \varphi'' + Sc(f\varphi' - \eta\varphi') - R\varphi = 0, \tag{7}$$

These are the associated boundary conditions

$$f(\eta) = 0, f''(\eta) = 0, f^{iv}(\eta) = 0, \theta(\eta) = 0, \varphi'(\eta) = 0 \text{ at } \eta = 0,$$

$$f(\eta) = 1, f'(\eta) = 0, \theta(\eta) = 1, \varphi(\eta) = 1 \text{ at } \eta = 1, \tag{14}$$

The following lists are the correlations between the thermophysical characteristics of the hybrid nanofluid:

$$\frac{\mu_{hnf}}{\mu_f} = \frac{1}{(1 - \phi_{hnf})^{2.5}} \text{ Dynamic viscosity } (\mu) \tag{15}$$

$$\rho_{hnf} = (1 - \phi_{hnf}) \rho_f + \phi_{V_2O_5} \rho_{V_2O_5} \text{ Density of hybrid nanofluid } (\rho) \tag{16}$$

$$(\rho Cp)_{hnf} = (1 - \phi_{hnf}) (\rho Cp)_f + \phi_{V_2O_5} (\rho Cp)_{V_2O_5} \text{ Heat capacity } (\rho Cp) \tag{17}$$

$$\frac{\sigma_{hnf}}{\sigma_f} = \left[1 + \frac{3\phi_{hnf} (\phi_{V_2O_5} \sigma_{V_2O_5})}{\phi_{V_2O_5} \sigma_{V_2O_5} + 2\phi_{hnf} \sigma_{hnf} - \phi_{hnf} \sigma_f} \right] \left(\phi_{V_2O_5} \sigma_{V_2O_5} - \sigma_f (\phi_{V_2O_5} \sigma_{V_2O_5}) \right) \tag{18}$$

Electrical conductivity (σ)

$$\frac{k_{hnf}}{k_f} = \frac{(\rho_{V_2O_5} k_{V_2O_5} + \phi_{hnf} k_{V_2O_5}) + 2k_f + (\phi_{V_2O_5} k_{V_2O_5} - 2\phi_{V_2O_5} k_f)}{\left(\frac{\phi_{V_2O_5} k_{V_2O_5}}{\phi_{V_2O_5}}\right) + 2k_f - (\phi_{V_2O_5} k_{V_2O_5}) + \phi_{V_2O_5} k_f}$$

Thermal conductivity (k) (19)

The substantial expressions from the governing equations are defined as

$$Rd = \frac{4\sigma^* T_\infty^3}{k^* k_f}, S = \frac{\alpha l^2}{2\nu_f}, Ha = lB_0 \sqrt{\frac{\sigma}{\rho_f \nu_f}}, R = \frac{ak_2 l^2}{\nu_f}$$

$$Sc = \frac{\nu_f}{D_m}, Pr = \frac{\nu_f}{\alpha_f}, Da = \frac{k_0}{\phi l^2}, \gamma = \frac{Q_0 l^2}{\nu_f (\rho Cp)_f}$$

$$Ec = \frac{\alpha^2 x^2}{4C_p T_w (1-at)^2}, De = \frac{\alpha \lambda_2}{1-at}, \delta = \frac{1}{x} (1-at)^{\frac{1}{2}}$$

4. Results and Discussion

This portion includes a brief explanation for better comprehension as well as graphical representations of the characteristics of fluid profiles for many relevant parameters. For hybrid nanofluid containing vanadium pentoxide, the important features of the heat radiation parameter, squeezing number, Eckert number, chemical reaction, ratio of retardation time parameter, and some other relevant factors on fluid profiles are described in depth. Additionally, table 1 above discusses the thermophysical characteristics of hybrid nanoparticles. By applying appropriate similarity transformations, the flow's governing equations are converted into ODEs. Using a shooting technique, via the Runge–Kutta–Fehlberg fourth–fifth-order (RKF 45) approach, the modified collection of ODEs is solved. To validate the current finding for $-f'(1)$, Table 2 has been produced using published results from Noor and Shafie [28] and Jyothi et al. [36]. Whereas, table 3, shows the computed values for the reduced Sherwood number concerning the thermal radiation parameter and the squeeze number. It is clear that irrespective of the values of the thermal radiation the result remains unchanged. No wonder because the radiation parameter does not appear in the concentration equation, hence, a change in its values may not necessarily make any significant change. For the values of $(S < 0)$ including $(S = 0)$, it can be heeded that an increase in the values of the parameter decreases the Sherwood number. In contrast, for the values of $(S > 0)$, it can be seen clearly that raise in the values of the parameter increases the values of the Sherwood number. The momentum profile $f(\eta)$ for a range of Deborah number De values is shown in Figure 2 depicts the thermodynamic changes in the temperature profile for Deborah number (De). It is observed that as Deborah's number increases, the thermal field decreases. The reason for this is that a rise in Deborah's number lowers the fluid's kinematic viscosity, which lowers the system's temperature. Furthermore, Figure 3 illustrates the temperature inclining with a higher Hartmann number. It is

demonstrated that the fluid's temperature increases when a magnetic field is applied to its lower surface. The flow's resistance is increased by the Lorentz force that MHD produces.

Table 2.0 Numerical results of $-f''(1)$ for Squeeze number (S) With $\lambda_1 \rightarrow \infty, Da \rightarrow \infty, De \rightarrow \infty, Ha = Ec = \delta = \gamma = R = \phi_2 = 0$ and $Sc = Pr = 1$.

Squeeze number (s)	Jyothi et al. [36] $-f''(1)$	Noor and Shafie [28] $-f''(1)$	Present result $-f''(1)$
-1.0	2.170090	2.170255	2.170223
-0.5	2.617403	2.617512	2.617391
0.01	3.007133	3.007208	3.007001
0.5	3.336449	3.336504	3.336620
2.0	4.167389	4.167411	4.167389

Table 3.0 Computations showing the reduced Sherwood number ($-\phi'(0)$) When

$Pr=15, \gamma=1.5, Ec=5.0, \delta=14, Sc=1, R=3$

Thermal radiation (Rd)	Sher $S=-5$	Sher $S=-1$	Sher $S=0$	Sher $S=1$	Sher $S=5$
0	0.8926	0.8827	0.6326	0.6619	0.7021
1	0.8926	0.8827	0.6326	0.6619	0.7021
3	0.8926	0.8827	0.6326	0.6619	0.7021
5	0.8926	0.8827	0.6326	0.6619	0.7021
7	0.8926	0.8827	0.6326	0.6619	0.7021

Figure 4 shows how the temperature profile is affected by the Eckert number Ec . As the Eckert number increases, so does the temperature profile. The reason for this is that heat energy is released into the fluid due to frictional forces, which increases the temperature field in the flow region. In addition, the temperature field rises when viscous dissipation occurs. Since the temperature equation directly mentions the Eckert number, adjusting the number's values can readily influence the temperature field.

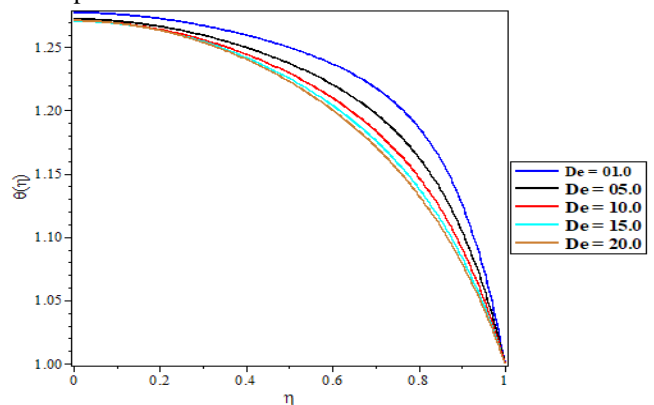


Fig. 2 Effects of Deborah Number on Temperature Profile

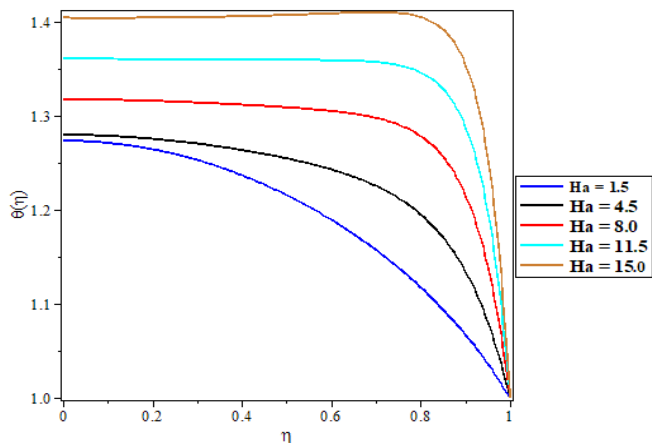


Fig. 3 Effects of Hartmann Number on Temperature Profile

The effects of the retardation time parameter (λ) ratio on the temperature field are shown in Figure 5. It can be seen that the temperature profile drops and the thermal field grows as (λ) increases. This is due to the working fluid's temperature rising during the heat generation process. The flow zone experiences an increase in the thermal field as a result. Moreover, an exothermic chemical reaction can lower the temperature field.

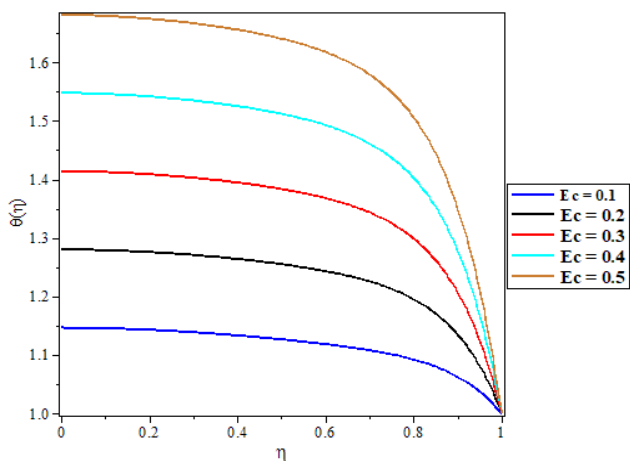


Fig. 4 Effects of Eckert Number on Temperature Profile

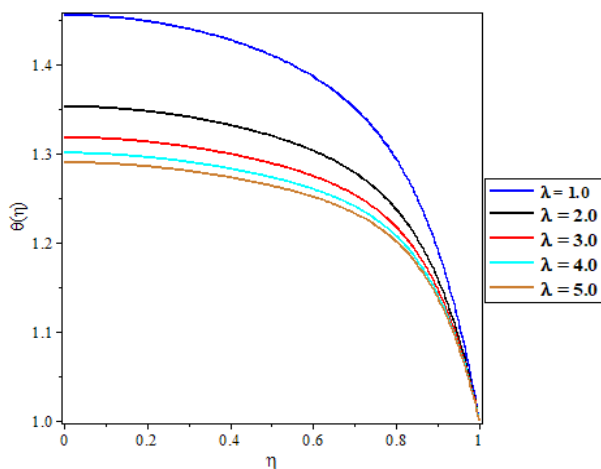


Fig. 5 Effects of Ratio of Retardation Time on Temperature Profile

Figure 6 shows the thermodynamic changes in flow profiles that are seen when the squeezing number (S) is applied. In the current investigation, the values of $S > 0$ indicate that the plates are moving apart from each other, while $S < 0$ indicates that they are getting closer. The relationship between squeeze number and temperature profile shows that for $S < 0$, the temperature profile drops, while for $S > 0$, it grows. The longer distance between the plates causes the fluid's kinematic viscosity to drop, which in turn increases the speed of the plates and lowers the temperature field. This is the cause of the temperature field's decline. Figure 7 is on the effects of thermal radiation Rd on the temperature profile, on this it can be observed that for any increment in the values of radiation parameter the temperature

Furthermore, the effects of thermal radiation on the Nusselt number profile are depicted Figure 8 makes it clear that a rise in the radiation parameter values results in a corresponding decrease in the Nusselt number profile. Figure 9 illustrates how the thickness of the thermal boundary layer reduces as the curves get steeper and the Prandtl number rises. The result is an increase in the lowered Nusselt number, which is proportionate to the starting slope. The free convective boundary layer flow in a typical fluid is similar to this pattern. Figure 10 portrays the concentration declining with an increase in the Hartmann number Ha . Empirical evidence indicates that the lower surface magnetic field slows down the fluid concentration.

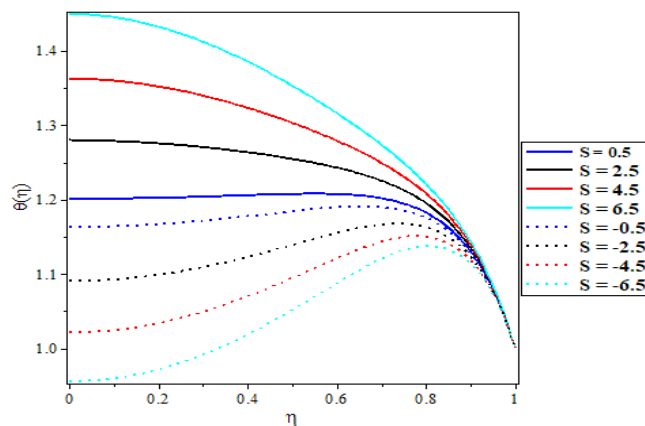


Fig. 6 Effects of Squeeze Number on Temperature Profile

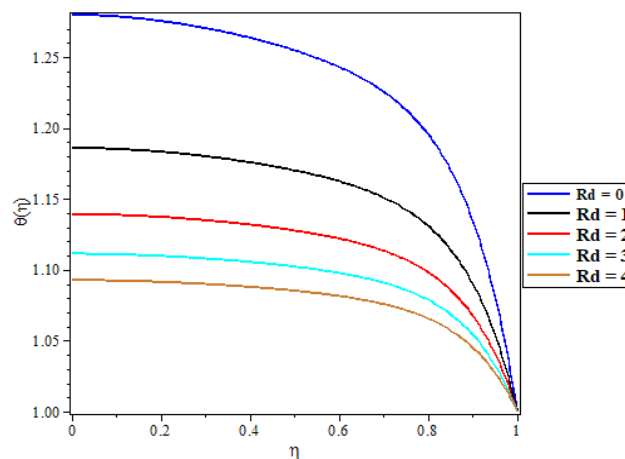


Fig. 7 Effects of Radiation parameter on Temperature Profile

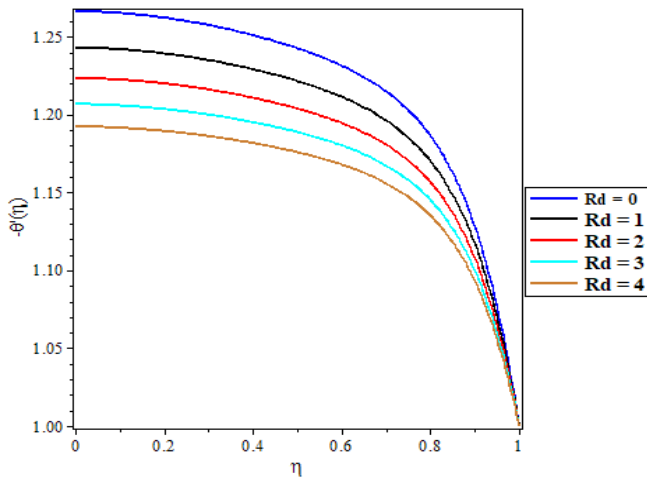


Fig. 8 Effects of Radiation Parameter on Nusselt Number Profile

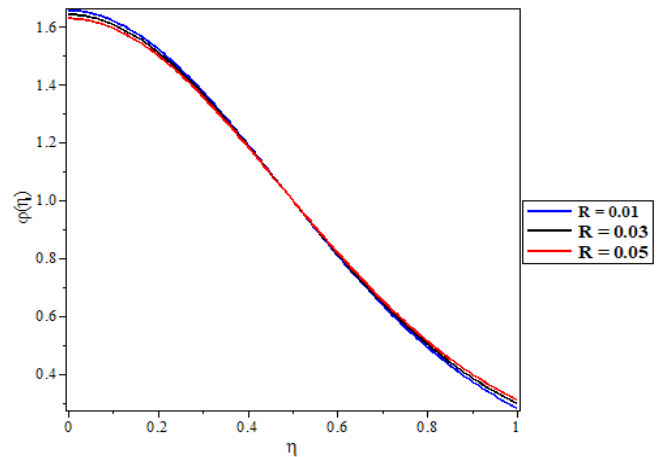


Fig. 11 Effects of Chemical Reaction Parameter on Concentration Profile

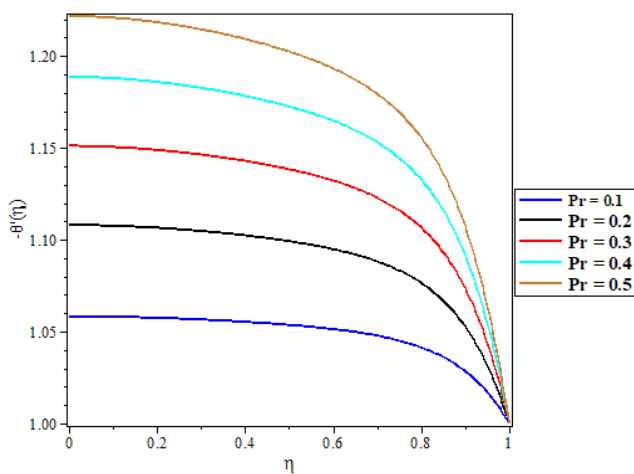


Fig. 9 Effects of Prandtl Number on Nusselt Number Profile

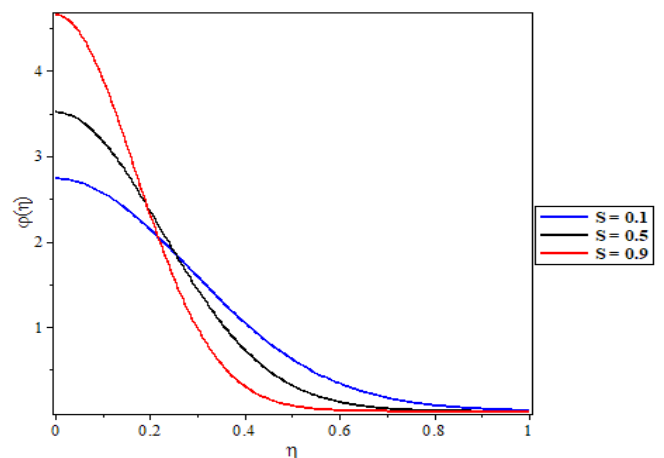


Fig. 12 Effects of Squeeze Number on Concentration Profile

The flow's resistance is increased by the Lorentz force that MHD produces. Figure 11 shows how the chemical reaction parameter R affects the Nusselt number profile; in this instance, an increase in the chemical reaction parameter results in a corresponding drop in the Nusselt number profile. And lastly, the effects of the squeeze number over the profile of concentration are displayed in Figure 12, in such a way that any increment in the squeeze number, produces an increment in the concentration profile.

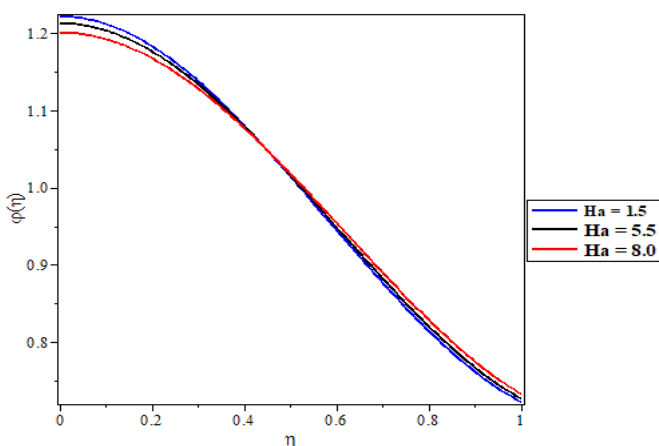


Fig. 10 Effects of Hartmann Number on Concentration Profile

5. Conclusion and Future Scope

For this study, we investigated the Casson hybrid nanofluid flow behaviour between two parallel plates while taking heat radiation and particle deposition of vanadium pentoxide (V_2O_5) into account. The governing equations that illustrate the fluid flow are then converted into non-linear ODEs with the support of appropriate similarity variables. The resulting equations are then numerically solved using the shooting technique and the RKF 45 scheme. Subsequently, Graphs are used to show how some relevant parameters affect various fluid profiles. The following is a summary of the findings above:

1. The momentum profile accelerates with an increase in Deborah's number, while the reverse is observed for the temperature profile.
2. The temperature and concentration profiles are enhanced with an increase in the Hartmann number.
3. The temperature is seen to increase with an increase in the Eckert number.
4. An increase in the ratio retardation time parameter, results in a decrease in the temperature, while in the case of concentration, reverse is observed.
5. The momentum profile is said to decrease as soon as the planes remain squeezed ($S > 0$) while it slows down when the planes are disjoined ($S < 0$) in the medium of the

channel. While in both cases the temperature profile is said to increase.

6. The temperature profile is seen to increase with an increase in thermal radiation parameter whereas the Nusselt number profile is seen to decrease.
7. The Nusselt number profile is increased with an increase in the Prandtl number
8. The concentration profile is decreased with an increase in the chemical reaction parameter.

Data Availability (Size 10 Bold)

None.

Conflict of Interest

The authors declare that they have no competing interests.

Funding Source

This study was supported by the TETFund Institution Based Research (IBR) [grant number ATBU/DVC/ACAD/12 (BATCH 21)].

Authors' Contributions

A.G. Madaki initiated the work, studied the literature and developed the study. A.A. Hussaini is concerned with protocol outcomes, achieving ethical endorsement, patient recruitment, and data analysis. Lare wrote the first draft of the manuscript and wrote the final draft of the manuscript. All authors reviewed and edited the manuscript and approved the final version of the manuscript.

Acknowledgements

The editors and reviewers deserve a great deal of gratitude from the authors for their expertise and guidance throughout our research. The authors are especially grateful for their insightful remarks, which have surely raised the calibre of this study.

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